

A Transformed Version of the Theorem on the Expansion of Multiple Stochastic Ito Integrals of Any Arbitrary Multiplicity k , Based on the Multiple Fourier Series Converging in the Mean-Square Sense

Dmitriy F. Kuznetsov

Peter the Great Saint-Petersburg Polytechnic University, Russia

E-mail: sde_kuznetsov@inbox.ru

Let $J[\psi^{(k)}]_{T,t}$ be a multiple Ito stochastic integral:

$$J[\psi^{(k)}]_{T,t} = \int_t^T \psi_k(t_k) \dots \int_t^{t_2} \psi_1(t_1) d\mathbf{w}_{t_1}^{(i_1)} \dots d\mathbf{w}_{t_k}^{(i_k)},$$

where every $\psi_l(\tau)$ ($l = 1, \dots, k$) is a continuous function on $[t, T]$; $\mathbf{w}_\tau^{(i)} = \mathbf{f}_\tau^{(i)}$ for $i = 1, \dots, m$ and $\mathbf{w}_\tau^{(0)} = \tau$; $\mathbf{f}_\tau$ is a standard m -dimensional Wiener stochastic process with independent components $\mathbf{f}_\tau^{(i)}$ ($i = 1, \dots, m$); $i_1, \dots, i_k = 0, 1, \dots, m$.

Define the following function on a hypercube $[t, T]^k$:

$$K(t_1, \dots, t_k) = \begin{cases} \psi_1(t_1) \dots \psi_k(t_k), & t_1 < \dots < t_k \\ 0, & \text{иначе} \end{cases} ; t_1, \dots, t_k \in [t, T]; k \geq 2,$$

and

$$K(t_1) = \psi_1(t_1); t_1 \in [t, T].$$

Suppose that $\{\phi_j(x)\}_{j=0}^\infty$ is a complete orthonormal system of functions in $L_2([t, T])$. The function $K(t_1, \dots, t_k)$ is sectionally continuous in the hypercube $[t, T]^k$. At this situation it is well known, that the multiple Fourier series of $K(t_1, \dots, t_k) \in L_2([t, T]^k)$ is converging to $K(t_1, \dots, t_k)$ in the hypercube $[t, T]^k$ in the mean-square sense, i.e.

$$\lim_{p_1, \dots, p_k \rightarrow \infty} \left\| K(t_1, \dots, t_k) - \sum_{j_1=0}^{p_1} \dots \sum_{j_k=0}^{p_k} C_{j_k \dots j_1} \prod_{l=1}^k \phi_{j_l}(t_l) \right\| = 0,$$

where

$$C_{j_k \dots j_1} = \int_{[t, T]^k} K(t_1, \dots, t_k) \prod_{l=1}^k \phi_{j_l}(t_l) dt_1 \dots dt_k, \quad (1)$$

$$\|f\|^2 = \int_{[t, T]^k} f^2(t_1, \dots, t_k) dt_1 \dots dt_k.$$

Consider the partition $\{\tau_j\}_{j=0}^N$ of $[t, T]$ such that

$$t = \tau_0 < \dots < \tau_N = T, \Delta_N = \max_{0 \leq j \leq N-1} \Delta\tau_j \rightarrow 0 \text{ if } N \rightarrow \infty, \Delta\tau_j = \tau_{j+1} - \tau_j. \quad (2)$$

Theorem 1 (see [1], [2] - [9]). Suppose that every $\psi_l(\tau)$ ($l = 1, \dots, k$) is a continuous on $[t, T]$ function and $\{\phi_j(x)\}_{j=0}^\infty$ is a complete orthonormal system of continuous functions in $L_2([t, T])$. Then

$$J[\psi^{(k)}]_{T,t} = \lim_{p_1, \dots, p_k \rightarrow \infty} \sum_{j_1=0}^{p_1} \dots \sum_{j_k=0}^{p_k} C_{j_k \dots j_1} \left(\prod_{l=1}^k \zeta_{j_l}^{(i_l)} - \lim_{N \rightarrow \infty} \sum_{(l_1, \dots, l_k) \in \mathcal{G}_k} \phi_{j_1}(\tau_{l_1}) \Delta \mathbf{w}_{\tau_{l_1}}^{(i_1)} \dots \phi_{j_k}(\tau_{l_k}) \Delta \mathbf{w}_{\tau_{l_k}}^{(i_k)} \right), \quad (3)$$

where

$$\mathcal{G}_k = \mathcal{H}_k \setminus \mathcal{L}_k; \quad \mathcal{H}_k = \{(l_1, \dots, l_k) : l_1, \dots, l_k = 0, 1, \dots, N-1\};$$

$$\mathcal{L}_k = \{(l_1, \dots, l_k) : l_1, \dots, l_k = 0, 1, \dots, N-1; l_g \neq l_r (g \neq r); g, r = 1, \dots, k\};$$

l.i.m. is a limit in the mean-square sense; $i_1, \dots, i_k = 0, 1, \dots, m$; every

$$\zeta_j^{(i)} = \int_t^T \phi_j(s) d\mathbf{w}_s^{(i)}$$

is a standard Gaussian random variable for various i or j (if $i \neq 0$); $C_{j_k \dots j_1}$ is the Fourier coefficient (1); $\Delta \mathbf{w}_{\tau_j}^{(i)} = \mathbf{w}_{\tau_{j+1}}^{(i)} - \mathbf{w}_{\tau_j}^{(i)}$ ($i = 0, 1, \dots, m$); $\{\tau_j\}_{j=0}^{N-1}$ is a partition of $[t, T]$, which satisfies the condition (2).

It was shown in [2] - [9] that theorem 1 is valid for convergence in the mean of degree $2n$, $n \in \mathbb{N}$.

In order to evaluate significance of the theorem 1 for practice we will demonstrate its transformed particular cases for $k = 1, \dots, 6$:

$$J[\psi^{(1)}]_{T,t} = \lim_{p_1 \rightarrow \infty} \sum_{j_1=0}^{p_1} C_{j_1} \zeta_{j_1}^{(i_1)}, \quad (4)$$

$$J[\psi^{(2)}]_{T,t} = \lim_{p_1, p_2 \rightarrow \infty} \sum_{j_1=0}^{p_1} \sum_{j_2=0}^{p_2} C_{j_2 j_1} \left(\zeta_{j_1}^{(i_1)} \zeta_{j_2}^{(i_2)} - \mathbf{1}_{\{i_1=i_2 \neq 0\}} \mathbf{1}_{\{j_1=j_2\}} \right), \quad (5)$$

$$J[\psi^{(3)}]_{T,t} = \lim_{p_1, \dots, p_3 \rightarrow \infty} \sum_{j_1=0}^{p_1} \sum_{j_2=0}^{p_2} \sum_{j_3=0}^{p_3} C_{j_3 j_2 j_1} \left(\zeta_{j_1}^{(i_1)} \zeta_{j_2}^{(i_2)} \zeta_{j_3}^{(i_3)} - \mathbf{1}_{\{i_1=i_2 \neq 0\}} \mathbf{1}_{\{j_1=j_2\}} \zeta_{j_3}^{(i_3)} - \mathbf{1}_{\{i_2=i_3 \neq 0\}} \mathbf{1}_{\{j_2=j_3\}} \zeta_{j_1}^{(i_1)} - \mathbf{1}_{\{i_1=i_3 \neq 0\}} \mathbf{1}_{\{j_1=j_3\}} \zeta_{j_2}^{(i_2)} \right), \quad (6)$$

$$\begin{aligned} J[\psi^{(4)}]_{T,t} = & \lim_{p_1, \dots, p_4 \rightarrow \infty} \sum_{j_1=0}^{p_1} \dots \sum_{j_4=0}^{p_4} C_{j_4 \dots j_1} \left(\prod_{l=1}^4 \zeta_{j_l}^{(i_l)} - \right. \\ & - \mathbf{1}_{\{i_1=i_2 \neq 0\}} \mathbf{1}_{\{j_1=j_2\}} \zeta_{j_3}^{(i_3)} \zeta_{j_4}^{(i_4)} - \mathbf{1}_{\{i_1=i_3 \neq 0\}} \mathbf{1}_{\{j_1=j_3\}} \zeta_{j_2}^{(i_2)} \zeta_{j_4}^{(i_4)} - \\ & - \mathbf{1}_{\{i_1=i_4 \neq 0\}} \mathbf{1}_{\{j_1=j_4\}} \zeta_{j_2}^{(i_2)} \zeta_{j_3}^{(i_3)} - \mathbf{1}_{\{i_2=i_3 \neq 0\}} \mathbf{1}_{\{j_2=j_3\}} \zeta_{j_1}^{(i_1)} \zeta_{j_4}^{(i_4)} - \\ & - \mathbf{1}_{\{i_2=i_4 \neq 0\}} \mathbf{1}_{\{j_2=j_4\}} \zeta_{j_1}^{(i_1)} \zeta_{j_3}^{(i_3)} - \mathbf{1}_{\{i_3=i_4 \neq 0\}} \mathbf{1}_{\{j_3=j_4\}} \zeta_{j_1}^{(i_1)} \zeta_{j_2}^{(i_2)} + \\ & + \mathbf{1}_{\{i_1=i_2 \neq 0\}} \mathbf{1}_{\{j_1=j_2\}} \mathbf{1}_{\{i_3=i_4 \neq 0\}} \mathbf{1}_{\{j_3=j_4\}} + \mathbf{1}_{\{i_1=i_3 \neq 0\}} \mathbf{1}_{\{j_1=j_3\}} \mathbf{1}_{\{i_2=i_4 \neq 0\}} \mathbf{1}_{\{j_2=j_4\}} + \\ & \left. + \mathbf{1}_{\{i_1=i_4 \neq 0\}} \mathbf{1}_{\{j_1=j_4\}} \mathbf{1}_{\{i_2=i_3 \neq 0\}} \mathbf{1}_{\{j_2=j_3\}} \right), \quad (7) \end{aligned}$$

$$\begin{aligned}
& -\mathbf{1}_{\{i_6=i_5 \neq 0\}} \mathbf{1}_{\{j_6=j_5\}} \mathbf{1}_{\{i_1=i_4 \neq 0\}} \mathbf{1}_{\{j_1=j_4\}} \mathbf{1}_{\{i_2=i_3 \neq 0\}} \mathbf{1}_{\{j_2=j_3\}} \\
& -\mathbf{1}_{\{i_6=i_5 \neq 0\}} \mathbf{1}_{\{j_6=j_5\}} \mathbf{1}_{\{i_1=i_2 \neq 0\}} \mathbf{1}_{\{j_1=j_2\}} \mathbf{1}_{\{i_3=i_4 \neq 0\}} \mathbf{1}_{\{j_3=j_4\}} \\
& -\mathbf{1}_{\{i_6=i_5 \neq 0\}} \mathbf{1}_{\{j_6=j_5\}} \mathbf{1}_{\{i_1=i_3 \neq 0\}} \mathbf{1}_{\{j_1=j_3\}} \mathbf{1}_{\{i_2=i_4 \neq 0\}} \mathbf{1}_{\{j_2=j_4\}} \Big), \tag{9}
\end{aligned}$$

where $\mathbf{1}_A$ is the indicator of the set A .

Note, that rightness of formulas (4) – (9) can be verified by the fact, that if $i_1 = \dots = i_6 = i = 1, \dots, m$ and $\psi_1(s), \dots, \psi_6(s) \equiv \psi(s)$, then we can deduce (see [1], [2] – [9]) the following well known classical equalities which are right with probability 1:

$$\begin{aligned}
J[\psi^{(1)}]_{T,t} &= \frac{1}{1!} \delta_{T,t}, \\
J[\psi^{(2)}]_{T,t} &= \frac{1}{2!} \left(\delta_{T,t}^2 - \Delta_{T,t} \right), \\
J[\psi^{(3)}]_{T,t} &= \frac{1}{3!} \left(\delta_{T,t}^3 - 3\delta_{T,t} \Delta_{T,t} \right), \\
J[\psi^{(4)}]_{T,t} &= \frac{1}{4!} \left(\delta_{T,t}^4 - 6\delta_{T,t}^2 \Delta_{T,t} + 3\Delta_{T,t}^2 \right), \\
J[\psi^{(5)}]_{T,t} &= \frac{1}{5!} \left(\delta_{T,t}^5 - 10\delta_{T,t}^3 \Delta_{T,t} + 15\delta_{T,t} \Delta_{T,t}^2 \right), \\
J[\psi^{(6)}]_{T,t} &= \frac{1}{6!} \left(\delta_{T,t}^6 - 15\delta_{T,t}^4 \Delta_{T,t} + 45\delta_{T,t}^2 \Delta_{T,t}^2 - 15\Delta_{T,t}^3 \right),
\end{aligned}$$

where

$$\delta_{T,t} = \int_t^T \psi(s) d\mathbf{f}_s^{(i)}, \quad \Delta_{T,t} = \int_t^T \psi^2(s) ds,$$

which can be independently obtained using the Ito formula and Hermite polynomials [10].

Let's generalize formulas (4) – (9) for the case of any arbitrary multiplicity of $J[\psi^{(k)}]_{T,t}$. In order to do it we will introduce several denotations.

Let's examine the unregulated set $\{1, 2, \dots, k\}$ and separate it up in two parts: the first part consists of r unordered pairs (sequence order of these pairs is also unimportant) and the second one — of the remains $k - 2r$ numbers.

So, we have:

$$\left(\underbrace{\{\{g_1, g_2\}, \dots, \{g_{2r-1}, g_{2r}\}\}}_{\text{part 1}}, \underbrace{\{q_1, \dots, q_{k-2r}\}}_{\text{part 2}} \right), \tag{10}$$

where $\{g_1, g_2, \dots, g_{2r-1}, g_{2r}, q_1, \dots, q_{k-2r}\} = \{1, 2, \dots, k\}$, curly braces mean unordered set, and the round braces mean ordered set.

Let's call (10) as partition and examine the sum using all possible partitions:

$$\sum_{\substack{(\{g_1, g_2\}, \dots, \{g_{2r-1}, g_{2r}\}, \{q_1, \dots, q_{k-2r}\}) \\ \{g_1, g_2, \dots, g_{2r-1}, g_{2r}, q_1, \dots, q_{k-2r}\} = \{1, 2, \dots, k\}}} a_{g_1 g_2, \dots, g_{2r-1} g_{2r}, q_1 \dots q_{k-2r}}. \tag{11}$$

Let's give an example of sums in the form (11):

$$\sum_{\substack{(\{g_1, g_2\}) \\ \{g_1, g_2\} = \{1, 2\}}} a_{g_1 g_2} = a_{12},$$

$$\begin{aligned}
& \sum_{\substack{(\{g_1, g_2\}, \{g_3, g_4\}) \\ \{g_1, g_2, g_3, g_4\} = \{1, 2, 3, 4\}}} a_{g_1 g_2 g_3 g_4} = a_{1234} + a_{1324} + a_{2314}, \\
& \sum_{\substack{(\{g_1, g_2\}, \{q_1, q_2\}) \\ \{g_1, g_2, q_1, q_2\} = \{1, 2, 3, 4\}}} a_{g_1 g_2, q_1 q_2} = a_{12,34} + a_{13,24} + a_{14,23} + a_{23,14} + a_{24,13} + a_{34,12}, \\
& \sum_{\substack{(\{g_1, g_2\}, \{q_1, q_2, q_3\}) \\ \{g_1, g_2, q_1, q_2, q_3\} = \{1, 2, 3, 4, 5\}}} a_{g_1 g_2, q_1 q_2 q_3} = a_{12,345} + a_{13,245} + a_{14,235} + a_{15,234} + a_{23,145} + a_{24,135} + \\
& \quad + a_{25,134} + a_{34,125} + a_{35,124} + a_{45,123}, \\
& \sum_{\substack{(\{g_1, g_2\}, \{g_3, g_4\}, \{q_1\}) \\ \{g_1, g_2, g_3, g_4, q_1\} = \{1, 2, 3, 4, 5\}}} a_{g_1 g_2, g_3 g_4, q_1} = a_{12,34,5} + a_{13,24,5} + a_{14,23,5} + a_{12,35,4} + a_{13,25,4} + a_{15,23,4} + \\
& \quad + a_{12,54,3} + a_{15,24,3} + a_{14,25,3} + a_{15,34,2} + a_{13,54,2} + a_{14,53,2} + a_{52,34,1} + a_{53,24,1} + a_{54,23,1}.
\end{aligned}$$

Now we can formulate the theorem 1 (formula (3)) using alternative and more comfortable form.

Theorem 2 (see [2] - [9]). *In conditions of the theorem 1 the following converging in the mean-square sense expansion is valid:*

$$\begin{aligned}
J[\psi^{(k)}]_{T,t} &= \text{l.i.m.}_{p_1, \dots, p_k \rightarrow \infty} \sum_{j_1=0}^{p_1} \dots \sum_{j_k=0}^{p_k} C_{j_k \dots j_1} \left(\prod_{l=1}^k \zeta_{j_l}^{(i_l)} + \right. \\
&+ \sum_{r=1}^{[k/2]} (-1)^r \cdot \sum_{\substack{(\{g_1, g_2\}, \dots, \{g_{2r-1}, g_{2r}\}, \{q_1, \dots, q_{k-2r}\}) \\ \{g_1, g_2, \dots, g_{2r-1}, g_{2r}, q_1, \dots, q_{k-2r}\} = \{1, 2, \dots, k\}}} \prod_{s=1}^r \mathbf{1}_{\{i_{g_{2s-1}} = i_{g_{2s}} \neq 0\}} \mathbf{1}_{\{j_{g_{2s-1}} = j_{g_{2s}}\}} \prod_{l=1}^{k-2r} \zeta_{j_{q_l}}^{(i_{q_l})} \Big). \quad (12)
\end{aligned}$$

In particular from (12) if $k = 5$ we obtain:

$$\begin{aligned}
J[\psi^{(5)}]_{T,t} &= \sum_{j_1, \dots, j_5=0}^{\infty} C_{j_5 \dots j_1} \left(\prod_{l=1}^5 \zeta_{j_l}^{(i_l)} - \sum_{\substack{(\{g_1, g_2\}, \{q_1, q_2, q_3\}) \\ \{g_1, g_2, q_1, q_2, q_3\} = \{1, 2, 3, 4, 5\}}} \mathbf{1}_{\{i_{g_1} = i_{g_2} \neq 0\}} \mathbf{1}_{\{j_{g_1} = j_{g_2}\}} \prod_{l=1}^3 \zeta_{j_{q_l}}^{(i_{q_l})} + \right. \\
&+ \sum_{\substack{(\{g_1, g_2\}, \{g_3, g_4\}, \{q_1\}) \\ \{g_1, g_2, g_3, g_4, q_1\} = \{1, 2, 3, 4, 5\}}} \mathbf{1}_{\{i_{g_1} = i_{g_2} \neq 0\}} \mathbf{1}_{\{j_{g_1} = j_{g_2}\}} \mathbf{1}_{\{i_{g_3} = i_{g_4} \neq 0\}} \mathbf{1}_{\{j_{g_3} = j_{g_4}\}} \zeta_{j_{q_1}}^{(i_{q_1})} \Big).
\end{aligned}$$

The last equality obviously agree with (8).

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